NUMERICAL SOLUTIONS OF THE NONLINEAR POROUS MEDIA EQUATION BASED ON HIGH EXPLORATION PARTICLE SWARM OPTIMIZATION AND MOVING LEAST SQUARES

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Abstract

In this study, a new numerical method based on the combination of High Exploration Particle Swarm Optimization (HEPSO) and Moving Least Squares (MLS) is introduced to solve nonlinear porous media equations. The MLS scheme is employed to describe an appropriate discretized function, and the penalty method is implemented to convert the constrained problem into an unconstrained one via satisfying the initial conditions. The identified objective function is minimized by the HEPSO to find the approximated nodal values for the nonlinear porous media equation. In order to illustrate the effectiveness of the HEPSO, the optimization trajectories are compared with those of a Standard Particle Swarm Optimization (SPSO) algorithm. Moreover, comparisons are made between the exact solution and the introduced strategy to expose the accuracy, effectiveness and simplicity of the proposed method.

Keywords: Moving least square, penalty method, porous media equation, numerical solution, high exploration particle swarm optimization.

1. Introduction

Nonlinear partial differential equations as mathematical models frequently appear in many branches of science, engineering and economy. Many models have no closed form solutions and, hence, there is a need to seek approximate solutions by means of numerical methods. To name just a few, İzgi and Çetin (2018) implemented semi-implicit split-step numerical methods for a class of nonlinear stochastic differential equations with non-Lipschitz drift terms. Firoozjaee and Yousefi (2018) utilized using Ritz approximation and proposed a numerical approach for fractional partial differential equations. Fakhar-Izadi and Dehghan (2018) suggested a numerical approach based on the fully spectral collocation method for nonlinear parabolic partial integro-differential equations. In order to solve hyperbolic partial differential equations, Doha et al. (2019) successfully applied a shifted Jacobi spectral-Galerkin method. Zhou and Zhang (2019) employed one-leg methods for nonlinear stiff fractional differential equations with Caputo derivatives. Esmaeilbeigi et al. (2019) developed a finite difference based numerical scheme to

analyze time-dependent stochastic partial differential equations using an especial type of the partition of unity collocation method. Li and Wang (2019) combined Galerkin approaches and finite element methods to numerically solve Caputo-type partial differential equations. Yang et al. (2019) presented numerical blow-up analysis of the linearly implicit Euler method for nonlinear parabolic integro-differential equations. Zenget al. (2019) exploited a novel numerical approach based on the discrete least squares collocation method for two-dimensional nonlinear time-dependent partial differential equations.

On the other hand, porous media equations form an important class of dynamical systems which often occurs in nonlinear heat and mass transfer, combustion theory and flows in porous media problems (Polyanin and Zaitsev, 2004). For instance, it describes unsteady heat transfer in a quiescent medium with the heat diffusivity as a power-law function of a temperature (Pamuk, 2005). Consequently, many attempts have been conducted by researchers to investigate the nonlinear porous media equation in the last decade. For instance, Antoniouk and Arnaudon (2014) developed the variational principle to present the solutions of the weighted porous media equation. Jensen et al. (2014) investigated the porous media equations and resistance coefficients for coastal structures. Duque et al. (2015) studied the moving mesh method to analyze the porous media equation with variable exponent. Campos and Soler (2016) investigated the qualitative behavior and traveling waves for flux-saturated porous media equations arising in optimal mass transportation. Grillo and Muratori (2016) examined smoothing effects for the porous media equation on Cartan-Hadamard manifolds. Röckner et al. (2018) inspected stochastic porous media equation on general measure spaces with increasing Lipschitz nonlinearities. Dareiotis et al. (2019) reconnoitered entropy solutions for stochastic porous media equations. Duan et al. (2019) probed an energetic variational approach to represent numerical solutions for the porous media equation.

In pursuance to this, we propose modifications in the existing methods to solve porous media equations based on the new heuristic method introduced by Mahmoodabadi et al. (2014 and 2018). To this end, the MLS scheme is employed to describe an appropriate discretized function, and the penalty method is implemented to convert the constrained problem into an unconstrained one via satisfying the initial conditions. The identified objective function is minimized by the HEPSO to find the approximated nodal values for the nonlinear porous media equation and results are compared with SPSO. From the obtained results, it is observed that the proposed new method is consistent, stable and accurate to solve the nonlinear porous media equation.

The rest of this paper is structured as follows. Section 2 states the governing nonlinear porous media differential equation. Section 3 briefly introduces the SPSO and HEPSO algorithms. The MLS scheme and discretization of the porous media equation are mentioned in Sections 4 and 5, respectively. The simulation results and the comparisons are discussed in Section 6. Section 7 concludes the results of this study.

2. Porous media differential equation

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Consider the nonlinear porous media differential equation as follows (Pamuk, 2005).

$$\frac{\partial\psi}{\partial\tau} = \frac{\partial}{\partial\zeta} \Big(\psi^r \frac{\partial\psi}{\partial\zeta}\Big),\tag{1}$$

where, r is a rational number, ζ represents position coordinate, τ denotes the time, and ψ signifies desired variable of the equation. Let us respect r = 1 in Eq. (1); therefore, the equation is rewritten as follows.

$$\frac{\partial\psi}{\partial\tau} = \frac{\partial}{\partial\zeta} \left(\psi \frac{\partial\psi}{\partial\zeta} \right),\tag{2}$$

with the initial condition of

$$\psi(\zeta, 0) = \psi_{\zeta, 0},\tag{3}$$

and boundary conditions of

$$\psi(0,\tau) = \psi_{0,\tau} \text{ and } \psi(1,\tau) = \psi_{1,\tau}.$$
 (4)

3. High exploration particle swarm optimization

The general form of the Standard PSO (SPSO) equations could be written as follows (Sibalija, 2019):

$$\vec{\phi}_i(t+1) = w \,\vec{\phi}_i(t) + \alpha \vec{r_1} \left(\vec{\psi}_{pbest_i}(t) - \vec{\psi}_i(t) \right) + \beta \vec{r_2} \left(\vec{\psi}_{gbest}(t) - \vec{\psi}_i(t) \right) \tag{5}$$

$$\vec{\psi}_i(t+1) = \vec{\psi}_i(t) + \vec{\phi}_i(t+1)$$
(6)

where, $\vec{\psi}_i(t)$ and $\vec{\phi}_i(t)$ describe the position and velocity of the particle *i* at certain iteration *t*, respectively. Coefficients α and β denote the tendency of the particle to the personal and social successes, respectively $\vec{r_1}$ and $\vec{r_2}$ are random vectors in interval [0,1], and *w* illustrates the effect of the particle velocity at the previous iteration on the next one. Furthermore, $\vec{\psi}_{pbest_i}(t)$ and $\vec{\psi}_{gbest}(t)$ are personal and global best positions of the particle *i* and the swarm at iteration *t*, respectively.

In the HEPSO algorithm, two operators are added to the PSO equations. The first one is based on the multi-crossover mechanism of the genetic algorithm and the other one is obtained via artificial bee colony.

Genetic algorithm operation: According to multiple-crossover idea, a new formulation for the velocity of a particle is proposed as follows:

$$\vec{\phi}_i(t+1) = \vec{r} \left(\frac{\beta}{2} \, \vec{\psi}_{gbest}(t) - \vec{\psi}_{pbest_i}(t) - \vec{\psi}_i(t)\right),\tag{7}$$

where, \vec{r} is a vector of random values between [0,1].

Artificial bee colony operation: The main operator of the artificial bee colony algorithm is utilized to determine the new position of particle i in dimension d as the following equation.

$$\psi_i^d(t+1) = \psi_i^d(t) + (2r-1)(\psi_i^d(t) - \psi_j^d(t)), \tag{8}$$

where, r is a random value between [0,1]. j is an integer random number between 1 and the maximum number of particles.

4. Moving least square

The moving least square approximation is formulized in this section. At first, the shape function $\chi(\mathbf{x}_i)$ for a selected point $\mathbf{x}_i = [\tau_i, \zeta_i]^T \in \Psi$ is defined by:

$$\chi(\mathbf{x}_i) = \omega^T(\mathbf{x}_i)\Gamma^{-1}(\mathbf{x})\mathbf{E}(\mathbf{x}) \text{ for } \mathbf{x} \in \Psi_d, \tag{9}$$

where, $\omega^T(\mathbf{x}_i) = [1, \mathbf{x}_i^1, \mathbf{x}_i^2, \mathbf{x}_i^3, \dots, \mathbf{x}_i^m]$ is the monomial basis vector of order *m*; and matrices $\Gamma(\mathbf{x})$ and $\mathbf{E}(\mathbf{x})$ are calculated by the following relations.

$$\Gamma(\mathbf{x}) = \Omega^T \Lambda \Omega \,, \tag{10}$$

$$\mathbf{E}(\mathbf{x}) = \mathbf{\Omega}^T \mathbf{\Lambda},\tag{11}$$

where, $\Omega = [\omega(\mathbf{x}_1), \omega(\mathbf{x}_2), \dots, \omega(\mathbf{x}_n)]^T$ is the monomial basis matrix; *n* refers to the number of nodes in the supporting domain Ψ_x of the global domain Ψ ; and Λ represents the diagonal weight function matrix whose *jth* diagonal member is calculated by the following equation.

$$\lambda_j(\mathbf{x}_i) = \begin{cases} 1 - 6(\frac{\bar{\rho}_{ij}}{\rho})^2 + 8(\frac{\bar{\rho}_{ij}}{\rho})^3 - 3(\frac{\bar{\rho}_{ij}}{\rho})^4 & 0 \le \bar{\rho}_{ij} \le \rho \\ 0 & \bar{\rho}_{ij} \ge \rho \end{cases},$$
(12)

where, j = 1, 2, ..., n; $\bar{\rho}_{ij}$ is the distance between points x_i and x_j ; and ρ represents the size of the supporting domain. After calculation of the shape function $\chi(x_i)$, the unknown value $\psi^h(x_i)$ is approximated from the nodal values $\hat{\psi}(x)$ as below (Mahmoodabadi et al., 2018)

$$\psi^{h}(\mathbf{x}_{i}) = \chi(\mathbf{x}_{i})\widehat{\psi}(\mathbf{x}). \tag{13}$$

5. Discretization of the porous media equation

In this section, the objective function of the problem is identified for minimization by the detailed HEPSO algorithm. Suppose that $x_i = [\tau_i, \zeta_i]^T$ denotes *i*th nodal position on the problem domain, and $\hat{\psi}(\mathbf{x}_i)$ indicates the approximation function value at \mathbf{x}_i . Using Eq. (13), Eq. (2) could be rewritten for *i*th node as follows.

$$\frac{d\chi(\mathbf{x}_i)}{d\tau}\hat{\psi}(\mathbf{x}_i) - \frac{d\chi(\mathbf{x}_i)}{d\zeta}[\hat{\psi}(\mathbf{x}_i)(\frac{d\chi(\mathbf{x}_i)}{d\zeta}\hat{\psi}(\mathbf{x}_i))] = 0$$
(14)

The summation of absolute values of these equations for all nodal points generates the first part of the discretized objective function.

$$\sum_{i=1}^{M} \left| \frac{d\chi(\mathbf{x}_{i})}{d\tau} \hat{\psi}(\mathbf{x}_{i}) - \frac{d\chi(\mathbf{x}_{i})}{d\zeta} [\hat{\psi}(\mathbf{x}_{i}) (\frac{d\chi(\mathbf{x}_{i})}{d\zeta} \hat{\psi}(\mathbf{x}_{i}))] \right| = 0,$$
(15)

where, M is the total number of nodes.

If the penalty method is utilized to impose the initial and boundary conditions, then:

$$\begin{split} \sum_{i=1}^{M} \left| \frac{d\chi(\mathbf{x}_{i})}{d\tau} \hat{\psi}(\mathbf{x}_{i}) - \frac{d\chi(\mathbf{x}_{i})}{d\zeta} [\hat{\psi}(\mathbf{x}_{i})(\frac{d\chi(\mathbf{x}_{i})}{d\zeta} \hat{\psi}(\mathbf{x}_{i}))] \right| + \Gamma_{1} |\psi(\zeta, 0) - \psi_{\zeta, 0}| + \Gamma_{2} |\psi(0, \tau) - \psi_{0, \tau}| + \Gamma_{3} |\psi(1, \tau) - \psi_{1, \tau}| &= 0 \end{split}$$

$$(16)$$

where, Γ_i (*i* = 1,2,3) \gg 1 are the penalty parameters.

6. Numerical results and comparisons

In this section, the introduced strategy based on the HEPSO and MLS is applied to analyze the porous media nonlinear differential equation illustrated by Eq. (2). If M is the total number of nodes, the population size of optimization algorithm is considered as $10 \times M$; its maximum number of iterations is set as $500 \times M$, while the penalty parameters are set at $\Gamma_i = 100$ (i = 1,2,3). The characteristics of the SPSO and HEPSO algorithms are presented in Table 1. The global domain is considered as $\tau \in [0,1]$ and $\zeta \in [0,1]$ with the supporting domain size $\rho =$

1.001 for different distances between nodal points in τ direction (h_{τ}) and ζ direction (h_{ζ}) . The problem is solved for several different initial and boundary conditions illustrated in the following examples.

6.1. Illustrative Example 1

Consider the problem of the porous media with governing Eq. (2) and the following initial and boundary conditions:

$$\psi(\zeta, 0) = \zeta, \psi(0, \tau) = \tau, \ \psi(1, \tau) = 1 + \tau \tag{17}$$

The exact solutions for this problem is as follows:

$$\psi(\zeta,\tau) = \zeta + \tau \tag{18}$$

Algorithm	Parameters			
SPSO	$\alpha = 3; \beta = 3; w = 1$			
HEPSO	$\alpha_{1i} = 2.5; \alpha_{1f} = 0.5; \beta_{2i} = 0.5; \beta_{2i} = 2.5; w = adaptive; P_C = 0.95; P_B = 0.02$			





Fig. 1. Convergence trajectory of SPSO and HEPSO for $h_{\tau} = 0.5$ and $h_{\zeta} = 0.5$ (Example 1).



Fig. 2. Convergence trajectory of SPSO and HEPSO for $h_{\tau} = 0.5$ and $h_{\zeta} = 0.25$ (Example 1).



Fig. 3. Convergence trajectory of SPSO and HEPSO for $h_{\tau} = 0.5$ and $h_{\zeta} = 0.2$ (Example 1).



Fig. 4. Convergence trajectory of SPSO and HEPSO for $h_{\tau} = 0.5$ and $h_{\zeta} = 0.125$ (Example 1).



Fig. 5. Convergence trajectory of SPSO and HEPSO for $h_{\tau} = 0.5$ and $h_{\zeta} = 0.1$ (Example 1).



Fig. 6. Convergence trajectory of SPSO and HEPSO for $h_{\tau} = 0.25$ and $h_{\zeta} = 0.5$ (Example 1).



Fig. 7. Convergence trajectory of SPSO and HEPSO for $h_{\tau} = 0.2$ and $h_{\zeta} = 0.5$ (Example 1).



Fig. 8. Convergence trajectory of SPSO and HEPSO for $h_{\tau} = 0.125$ and $h_{\zeta} = 0.5$ (Example 1).



Fig. 9. Convergence trajectory of SPSO and HEPSO for $h_{\tau} = 0.1$ and $h_{\zeta} = 0.5$ (Example 1).

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$h_{ au}$	h_{ζ}	Objective function of	Objective function	Error of	Error of
	,	HEPSO	of SPSO	HEPSO	SPSO
0.5	0.1	0.0060136	9.1232	9.4037e-05	0.07413
	0.125	0.0032618	6.2153	5.7964e-05	0.08864
	0.2	0.00036714	2.7687	8.8421e-06	0.040681
	0.25	0.00032032	2.1042	7.0401e-06	0.014804
	0.5	0.005225	0.19557	3.4672e-05	0.00055015
0.1	0.5	0.013239	1.4875	4.1813e-05	0.0012832
0.125		0.011228	0.65529	4.3519e-05	0.00056184
0.2		0.008176	0.2854	4.2103e-05	0.00044996
0.25		0.007133	0.23361	3.9744e-05	0.00036228

Table 2. Objective functions and errors for different distances between nodal points in directions τ and ζ (Example 1).



Fig. 10. Comparison of results obtained via the proposed method and exact solution for $h_{\tau} = 0.5$ and $h_{\zeta} = 0.5$ (Example 1).

Fig. 11. Comparison of results obtained via the proposed method and exact solution for $h_{\tau} = 0.5$ and $h_{\zeta} = 0.25$ (Example 1).

Fig. 12. Comparison of results obtained via the proposed method and exact solution for $h_{\tau} = 0.5$ and $h_{\zeta} = 0.2$ (Example 1).

Fig. 13. Comparison of results obtained via the proposed method and exact solution for $h_{\tau} = 0.5$ and $h_{\zeta} = 0.125$ (Example 1).

Fig. 14. Comparison of results obtained via the proposed method and exact solution for $h_{\tau} = 0.5$ and $h_{\zeta} = 0.1$ (Example 1).

Fig. 15. Comparison of results obtained via the proposed method and exact solution for $h_{\tau} = 0.25$ and $h_{\zeta} = 0.5$ (Example 1).

Fig. 16. Comparison of results obtained via the proposed method and exact solution for $h_{\tau} = 0.2$ and $h_{\zeta} = 0.5$ (Example 1).

Fig. 17. Comparison of results obtained via the proposed method and exact solution for $h_{\tau} = 0.125$ and $h_{\zeta} = 0.5$ (Example 1).

Fig. 18. Comparison of results obtained via the proposed method and exact solution for $h_{\tau} = 0.1$ and $h_{\zeta} = 0.5$ (Example 1).

The convergence trajectories of the SPSO and HEPSO are shown in Figs. 1 through 9 for different distances between nodal points (for more visibility, only the first 2000 iterations are exhibited). Further, the objective functions obtained by SPSO and HEPSO for the global best particle as well as the related errors are illustrated in Table 2. This table clearly depicts that the HEPSO approach is able to find better results in comparison with the SPSO method. Additionally, the numerical solutions found via the MLS based PSO scenario are compared with those of the

exact solution in Figs. 10 to 18 for different nodal distributions. These comparisons obviously portray the high accuracy, great stability and excellent convergence of the proposed strategy to solve the porous media nonlinear differential equation.

6.2. Illustrative Example 2

The governing porous media equation was considered as Eq. (2) and the initial and boundary conditions are regarded as follows.

$$\psi(\zeta, 0) = -\frac{\zeta^2}{6}, \psi(0, \tau) = 0, \ \psi(1, \tau) = -\frac{1}{6(1+\tau)}$$
(19)

The exact solution related to the above equation could be given as Eq. (20).

$$\psi(\zeta,\tau) = -\frac{\zeta^2}{6(1+\tau)} \tag{20}$$

The convergence trajectories obtained by the regarded optimization algorithms (SPSO and HEPSO) are indicated in Figs. 19 through 27 for different number of the nodal points (for more visibility, only the first 5000 iterations are disclosed). Additionally, the final value of objective function of the global best particle and the related errors are represented in Table 3. This table clearly depicts that the HEPSO approach is able to find better results in comparison with the SPSO method. As well, the numerical solutions originated via the MLS based PSO method and the analytical approach are demonstrated in Figs. 28 to 36 for the various nodal distributions.

Fig. 19. Convergence trajectory of SPSO and HEPSO for $h_{\tau} = 0.25$ and $h_{\zeta} = 0.25$ (Example 2).

Fig. 20. Convergence trajectory of SPSO and HEPSO for $h_{\tau} = 0.25$ and $h_{\zeta} = 0.2$ (Example 2).

Fig. 21. Convergence trajectory of SPSO and HEPSO for $h_{\tau} = 0.25$ and $h_{\zeta} = 0.125$ (Example 2).

Fig. 22. Convergence trajectory of SPSO and HEPSO for $h_{\tau} = 0.25$ and $h_{\zeta} = 0.1$ (Example 2).

Fig. 23. Convergence trajectory of SPSO and HEPSO for $h_{\tau} = 0.2$ and $h_{\zeta} = 0.25$ (Example 2).

Fig. 24. Convergence trajectory of SPSO and HEPSO for $h_{\tau} = 0.125$ and $h_{\zeta} = 0.25$ (Example 2).

Fig. 25. Convergence trajectory of SPSO and HEPSO for $h_{\tau} = 0.1$ and $h_{\zeta} = 0.25$ (Example 2).

$h_{ au}$	h_{ζ}	Objective function	Objective function	Error of	Error of
		of HEPSO	of SPSO	HEPSO	SPSO
0.25	0.1	0.056981	3.7632	0.00082616	0.03264
	0.125	0.015525	3.2015	0.00014218	0.030369
	0.2	0.010368	2.4376	7.6648e-05	0.024178
	0.25	0.0095065	2.2788	7.9057e-05	0.026548
0.1		0.089099	4.6552	0.0042164	0.03105
0.125	0.25	0.018598	3.574	0.00082393	0.029527
0.2		0.008846	1.6514	8.4606e-05	0.023025

Table 3. Objective functions and errors for different distances between nodal points in directions τ and ζ (Example 2).

Fig. 26. Comparison of results obtained via the proposed method and exact solution for $h_{\tau} = 0.25$ and $h_{\zeta} = 0.25$ (Example 2).

Fig. 27. Comparison of results obtained via the proposed method and exact solution for $h_{\tau} = 0.25$ and $h_{\zeta} = 0.2$ (Example 2).

Fig. 28. Comparison of results obtained via the proposed method and exact solution for $h_{\tau} = 0.25$ and $h_{\zeta} = 0.125$ (Example 2).

Fig. 29. Comparison of results obtained via the proposed method and exact solution for $h_{\tau} = 0.25$ and $h_{\zeta} = 0.1$ (Example 2).

Fig. 30. Comparison of results obtained via the proposed method and exact solution for $h_{\tau} = 0.2$ and $h_{\zeta} = 0.25$ (Example 2).

Fig. 31. Comparison of results obtained via the proposed method and exact solution for $h_{\tau} = 0.125$ and $h_{\zeta} = 0.25$ (Example 2).

7. Conclusions

In this study, a new combination of the optimization algorithms and the interpolation methods was described to analyze the nonlinear porous media differential equation. The moving least

square was employed to produce the discretized function, and the penalty method was implemented to satisfy the initial and boundary conditions. Convergence trajectories and global optimum positions evidently articulated that the considered optimization algorithm, namely HEPSO, displays a higher performance and accuracy in comparison with the SPSO. Further, the solutions of the nonlinear porous media differential equation via the proposed strategy were compared with those of the analytical formulation for various number of regular nodes. These results illustrated that the introduced method is able to solve this class of nonlinear differential equations with an acceptable convergence speed and high accuracy without any restrictive assumptions.

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